

Quasiperiodic operators with monotone potentials

I. Kachkovskiy, MSU

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Euler Institute, St. Petersburg

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Lecture 2: 1d monotone potentials

(plus the remainder from Lecture 1)

<https://users.math.msu.edu/users/ikachkov/minicourse.html>

$$(H(x)\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + f(x+nd)\Psi_n \quad \ell^2(\mathbb{Z})$$

$$H_n(x) = \begin{pmatrix} f(x) & 1 & & & & \\ & 1 & f(x+d) & & & \\ & & 1 & f(x+2d) & & \\ & & & \ddots & \ddots & \\ & 0 & & & & 1 \\ & & 0 & & & & 1 & f(x+(n-1)d) \end{pmatrix}$$

Let $N_n(x, E) = \# \{ (H_n(x)) \cap (-\infty, E] \}$ (counting function).

$$N(E) := \lim_{N \rightarrow \infty} \frac{1}{N} \int_0^1 N_n(x, E) dx$$

↑
integrated density of states = distribution function of the density of states measure.

$$dN = \int_0^1 \langle d|E\rangle_{H(x)} e_0, e_0 \rangle dx: \text{ spectral measure averaged in } x$$

$$M_n(x, E) = \begin{pmatrix} P_n(x, E) & -P_{n-1}(x+d, E) \\ P_{n-1}(x, E) & -P_{n-2}(x+d, E) \end{pmatrix} \quad \begin{array}{l} n\text{-step transfer matrix} \\ H(x)\Psi = E\Psi \end{array}$$

$$L(E) = \lim_{n \rightarrow \infty} \frac{1}{n} \int_0^1 \ln \|M_n(x, E)\| dx = \inf_n \frac{1}{n} \int_0^1 \ln \|M_n(x, E)\| dx$$

Thouless formula:

$$L(E) = \int_{\mathbb{R}} \ln |E - E'| dN(E')$$

Heuristics: $L(E) \approx \frac{1}{h} \int_0^1 \ln |P_h(x, E)| dx =$
 $= \frac{1}{h} \int_0^1 \sum_{j=0}^{n-1} \ln |E - E_j(x)| dx = \int_0^1 \ln |E - E'| N_n(dE')$
↑
counting measure

This can be made rigorous by considering $E \in \mathbb{C}^+$ and showing that L is subharmonic on \mathbb{C}^+ .

Let $(H\psi)_n = \psi_{n+1} + \psi_{n-1} + V_n \psi_n$ (general operator)

We say that ψ is a generalized eigenfunction if

- $|\psi(n)| \leq C_1 (1+|n|)^{C_2}$ (polynomially bounded)
- $H\psi = E\psi$ for some $E \in \mathbb{R}$
↑
generalized eigenvalue.

In general, $\psi \notin \ell^2(\mathbb{Z})$.

Theorem (Snoel') | The spectral measure of H is supported on the set of generalized eigenvalues of H .

Many modern proofs of AL use Green's function method.

$$\text{Let } H_N = H|_{[0, N-1]}, \quad H\psi = E\psi.$$

$$\text{Then } H_N \psi = E\psi - \psi_{-1}e_0 - \psi_N e_N$$

$$\psi_m = -(H_N - E)_{0m}^{-1} \psi_{-1} - (H_N - E)_{m, N-1}^{-1} \psi_N$$

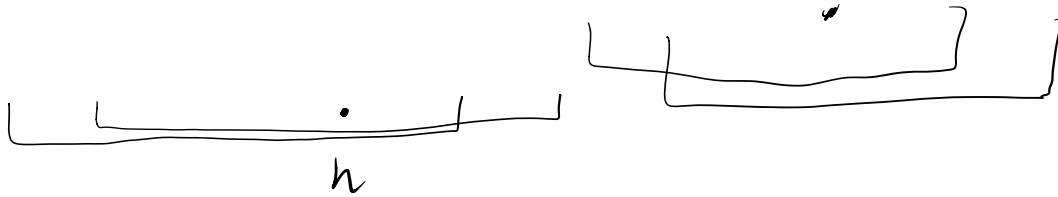
Then, try to prove that $(H_N - E)^{-1}$ has exponential decay away from the diagonal.

Cramer's rule:

$$(H_N - E)_{m, N-1}^{-1} = \frac{P_m(x, E)}{P_N(x, E)} \stackrel{?}{\sim} e^{(m-N)L(E)}$$

For most x , $P_m(x, E) \sim e^{mL(E)}$, $P_N(x, E) \sim e^{NL(E)}$
 In both cases, " \approx ". Therefore, need to avoid

points where $P_N(x, E) \ll e^{NL(E)}$.



Some freedom in choosing the box!

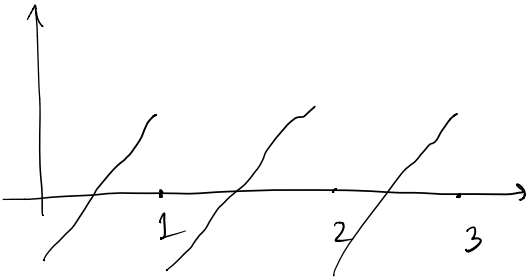
[Titomirskaya, 1998]

Relies on $L(E) > 0$.

$$(H(x) \Psi)_n = \Psi_{n+1} + \Psi_{n-1} + f(x+nd) \Psi_n$$

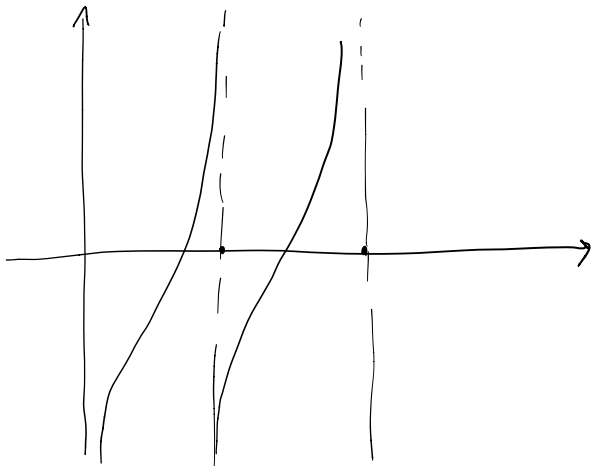
$f(x)$ is Lipschitz monotone on $(0, 1)$:

$$f(y) - f(x) \geq \gamma(y-x), \quad \gamma > 0, \quad 0 < x < y < 1.$$



$$f(x+d) = f(x)$$

← $\{x\}$ -type (bounded)



← Maryland-type (unbounded)

Example: $f(x) = \tan(\pi x)$

In both cases:

1) $L(E) > 0$ for Lebesgue a.e. E

2) $L(E) \geq \min\{\ln(\gamma/2e), 0\}$.

$$H_n(x) = \begin{pmatrix} f(x) & 1 & & & & \\ & 1 & f(x+d) & 1 & & \\ & & 1 & f(x+2d) & & \\ & & & \ddots & \ddots & \\ & 0 & & & 1 & \\ & & & & & 1 & f(x+(n-1)d) \end{pmatrix}$$

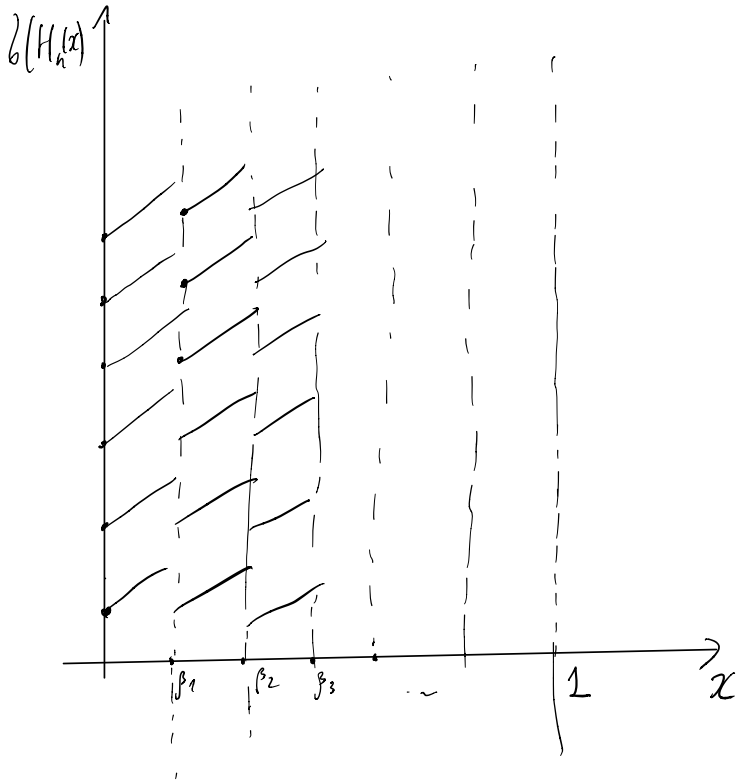
$$x \mapsto x + nd$$

$$[m, m+n]$$

Th In both cases, $|N(E_1) - N(E_2)| \leq \gamma^{-1} |E_1 - E_2|$
 IDS

Recall: need to look at counting functions:

$$\frac{1}{n} \int_0^1 N_n(x, E) dx, \quad N_n(x, E) = \# \delta(H_n(x)) \cap (-\infty, E]$$



$$\begin{pmatrix} f(x) \uparrow \\ 1 & f(x+2) \uparrow \\ & 1 & f(x+2\alpha) \uparrow \\ & & 1 & \dots \\ & & & & \dots \\ & & & & & & 1 \\ & & & & & & & 1 & f(x+(n-1)\alpha) \end{pmatrix} \begin{matrix} \circ \\ \circ \end{matrix}$$

$\{x\}$ -type case: ↗

Maryland-type case: ↘

$$\text{Let } Z = \{E : L(E) = 0\}$$

$$|Z| = 0 \implies N(Z) = 0 \implies$$

\nearrow
 density of states
 measure

$$\implies \int_0^1 \langle E_{H(x)}(z) e_j, e_j \rangle dx = 0.$$

Therefore, for a.e. x , Z has zero spectral measure!

Now we can restrict ourselves to $H(x)\Psi = E\Psi$, $L(E) = 0$.

$$\text{Let } P_n(x, E) = \det(H_n(x) - E) = \prod_{j=0}^{n-1} (E_j(x) - E)$$

The "jump points" on the above picture are at $\{ -jd \}$, $j = 0, 1, \dots, n-1$.

Let $\frac{p_j}{q_j} \rightarrow d$ (continued fractions), and

let q be one of the q_j . We will consider $n=q$.

The jump points are distributed regularly:

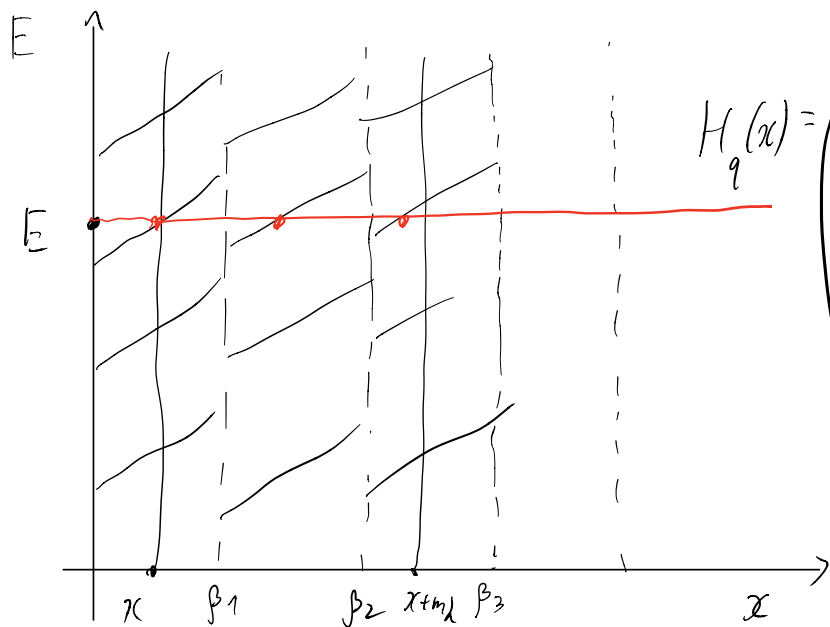
$$\beta_0 \quad \beta_1 \quad \beta_2 \quad \dots \quad \beta_{q-1}$$

$$\frac{1}{2q} \leq |\beta_{j+1} - \beta_j| \leq \frac{3}{2q}$$

Each interval $(\frac{j}{q}, \frac{j+1}{q})$ contains one β_j

$$H_q(x) \approx H_q(x+m d)$$

counting functions differ by a constant



$$H_q(x) = \begin{pmatrix} f(x) & 1 & & & \\ & 1 & f(x+d) & & 0 \\ & & 1 & \ddots & \\ 0 & & & & 1 \\ & 0 & & & & 1 & f(x+(q-1)d) \end{pmatrix}$$

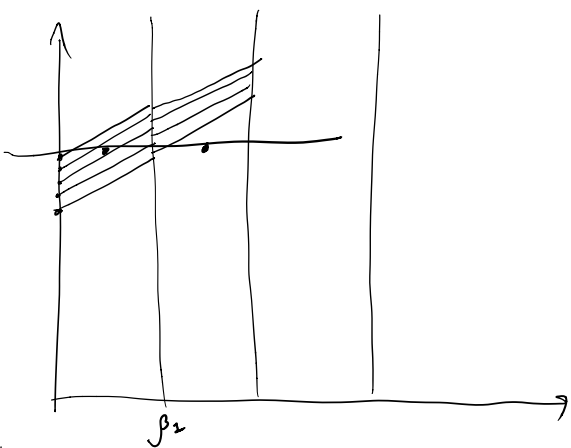
$$x \mapsto x + m d.$$

$$\{q d\}$$

$$P_n(x, E) = \det(H_n(x) - E)$$

Ultimately, $|N_q(x, E) - N_q(y, E)| \leq C$

Separation of eigenvalues modulo
finite rank is as good as
 γ . (separation of $\{x + jd\}$)



$$j = 0, 1, \dots, q-1$$

Large deviation set: q intervals of size $O(e^{-cq})$

Call $m \in \mathbb{Z}$ q -singular if one cannot
find a good q -interval around m .

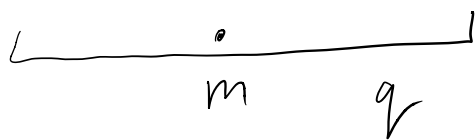
WLOG, 0 is q -singular for large q .

Therefore, $[q, e^{cq}]$ cannot contain singular points.

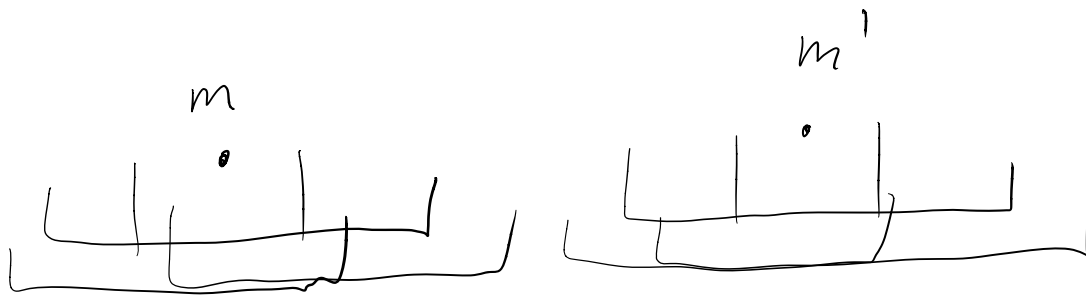
$$\Psi_m = - (H_N - E)_{0m}^{-1} \Psi_{-1} - (H_N - E)_{m, N-1}^{-1} \Psi_N$$

$$\frac{P_m(x, E)}{P_N(x, E)}$$

$$P_N(x, E)$$



q -regular



$$|m - m'| \geq \frac{q+1}{2}$$

both singular.

$P_N(x + m_j d)$ is small

$$2q(1 - o(1))$$

$m_j d$ in large deviation set

union of q intervals.

$$\text{dist}(m_j d - m_k d, \mathbb{Z}) < e^{-c q}$$

$$\begin{aligned} & \forall \\ & c |m_j - m_k|^{-\tau} \\ & |m_j - m_k| \gg e^{c_1 q} \end{aligned}$$

Theorem: Let f be γ -monotone

d be Diophantine. $\|nd\| \geq c|n|^{-\tau}$

Then $H(x)$ has localization:

1) for a.e. x

2) for all x if $\gamma > 2e$

Appendix 1:

Let V be an ergodic potential. We say that V is deterministic if $V|_{\mathbb{Z}_-}$ determines V .

Example: random potentials are not deterministic.

Th (Kotani). Let V be non-deterministic. Then $L(E) > 0$ for Lebesgue a.e. E .

$\{x\}$ seems deterministic. However, one can consider a "closure":

$$\begin{array}{ccc} \mathbb{T}^1 & \longrightarrow & [a, b]^{\mathbb{Z}} \\ \cup & & \cup \\ x & \longmapsto & \{x + n\omega\} : n \in \mathbb{Z} \\ d\theta & \longmapsto & d\mathcal{D}. \end{array}$$

(Haar measure)

The "extended" ergodic potential is not deterministic on $\text{supp } d\mathcal{D}$

(Damanik, Killip).

Unbounded case:

Th (Simon, Spencer) | Let V be unbounded. Then $\mathcal{B}_{ac}(H+V) = \emptyset$.

Therefore, in either case $L(E) > 0$ for a.e. E (Lebesgue), and we do not need to worry about δ_{ac} .
 However, the set $Z = \{E : L(E) = 0\}$ can still have positive spectral measure.

Appendix 2:

Th (Farman) Let (Ω, μ, T) be a uniquely ergodic system,

$f: \Omega \rightarrow [-\infty, +\infty)$ be subadditive:

$$f_{n+m}(x) \leq f_n(x) + f_m(T^n x),$$

$$f \in L^1(X, \mu).$$

$$\text{Let } L(f) = (\text{a.e.}) \lim_{n \rightarrow \infty} \frac{1}{n} f_n(x).$$

Then, uniformly on $x \in X$,

$$\lim_{n \rightarrow \infty} \sup \frac{1}{n} f_n(x) \leq L(f).$$

For $\{x\}$ -type, can apply directly.

For Maryland-type, "factor out" the unbounded part and analyze separately.

Appendix 3: Originally, this method was developed for the almost Mathieu operator:

$$(H\Psi)_n = \Psi_{n+1} + \Psi_{n-1} + 2\lambda \cos 2\pi(x + nd) \Psi_n.$$

In this case, P is a polynomial in $\cos x$.

n roots $\rightarrow 2n$ roots \rightarrow resonance.

In monotone case, no such resonances!

Leads to uniform localization

$$\Psi_n^{(j)} \leq C(\delta) e^{-(L(E)-\delta)|n-n_j|} \Psi_n^{(j)}(n_j)$$