

Symmetric Lévy processes with reflection

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Overview

Reflecting processes

Geometric reflection – c'est pas ça

Reflection is Neumann conditions for the generator

Local time huddles the process back into the domain

Example of Skorokhod reflection

Construction

A class of Lévy processes under consideration

Main idea: two extensions

Difference of semigroups lives on the boundary

Group structure of our operator families

Pathwise accumulated momentum

Limit theorem for the reflecting BM in d -ball

Part 1. General notions, some intuition and historical notes

Geometric reflection – c'est pas ça

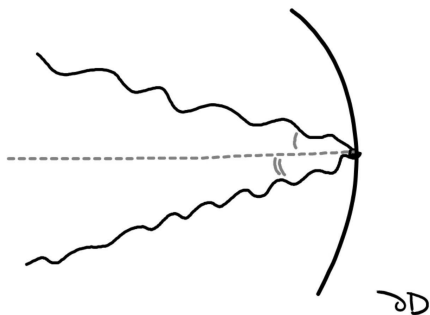
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Let $\xi(t)$ be a Markov process in \mathbb{R}^d . We shall call it *free* in opposition to its *confined to D* versions. We are interested in one single type of containment: *specular reflection* off the boundary. Note that we **cannot** define the reflection geometrically



since it means imposing conditions on normal derivatives, which very seldom exist.

Reflection is Neumann conditions for the generator

Taking average of f over the paths of ξ

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is a Markov contraction semigroup! We *define* the reflecting version of ξ by imposing Neumann conditions on its generator. Or on its quadratic form.

Example (see. [1])

Brownian Motion with reflection off the boundary ∂D is a process, associated with the Dirichlet form

$$\mathcal{E}(u, v) = \frac{1}{2} \int_D \nabla u \cdot \overline{\nabla v} \, d\mathbf{x}, \quad D[\mathcal{E}] = W_2^1(D).$$

Remark. Same form on $W_2^{1,0}(D)$ gives rise to an *absorbing* BM.

Local time huddles the process back into the domain

Skorokhod '61: $d = 1$

Skorokhod in [6] has proposed a *pathwise* construction of the reflecting diffusion in $[0, \infty)$ and proved the since-called Skorokhod semimartingale decomposition

$$\tilde{\xi}(t) = \xi(t) + \zeta(t),$$

where ζ is the local time of ξ .

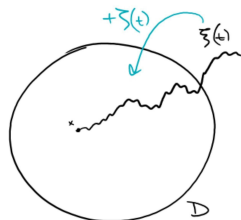


Figure: Local time pushes the process back into the domain D

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Case $d > 1$: $d\zeta$ is a measure on a boundary

If the boundary is C^3 -smooth, the Skorokhod construction is still applicable (see [3]), but in place of $\zeta(t)$ the decomposition features an **integral with respect to $d\zeta(t)$** .

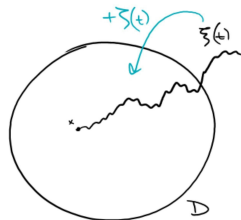
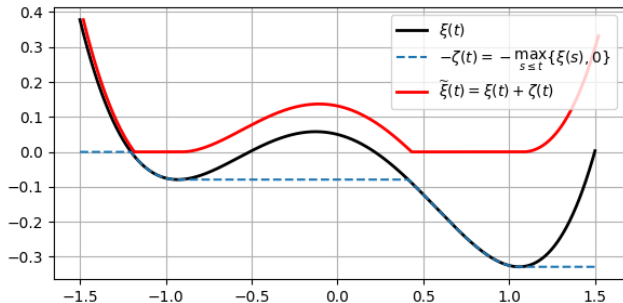


Figure: Local time pushes the process back into the domain D

Skorokhod reflection bears no resemblance to the geometric reflection

Let ξ be a *non-random* smooth function. Its reflecting version in $[0, \infty)$:



Here $\tilde{\xi}$ is the reflecting version and ζ is the compensator.

Part 2. Construction

A class of Lévy processes under consideration

consists of

1. pure jump processes (without diffusion terms)
2. with rotation-invariant Lévy measure $d\Pi$
3. and finite second moments

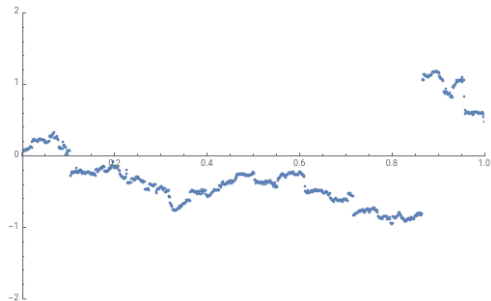


Figure: Typical path of α -stable process ($\alpha = 1.42$),
<https://demonstrations.wolfram.com/StableLevyProcess>

Characteristic function and generator

The characteristic function

is of the form $\varphi_t(\mathbf{p}) = \exp(-tL(\mathbf{p}))$, where

$$L(\mathbf{p}) = - \int_{\mathbb{R}^d} (e^{i\mathbf{p} \cdot \mathbf{x}} - 1 - i\mathbf{p} \cdot \mathbf{x}) d\Pi(\mathbf{x}). \quad (1)$$

Generator

$$-Lf(\mathbf{x}) = \int_{\mathbb{R}^d} (f(\mathbf{x} + \mathbf{y}) - f(\mathbf{x}) - f'(\mathbf{x}) \cdot \mathbf{y}) d\Pi(\mathbf{y}) \quad (2)$$

Remark. It *would* be a PDO if $L(\mathbf{p})$ was smooth (it in general is not!).

Main idea: two extensions

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1. take a function f in D
2. and extend it in some special way to the whole \mathbb{R}^d :
 - 2.1 \tilde{f} for the reflecting process
 - 2.2 \bar{f} for the process inside the domain
3. associate a semigroup to each by

$$P^t f = T^t \tilde{f}, \quad R^t f = T^t \bar{f},$$

denote its generators by $-A$ and $-A^N$.

Adjustment to the main idea

The idea works well when D is very simple (such as a ball or a cube). In an arbitrary smooth domain we construct instead two sequences $\tilde{f}_M(\mathbf{x}, \mathbf{y})$ and $\bar{f}_M(\mathbf{x}, \mathbf{y})$, which are **tangent to f** at $\mathbf{y} = 0$.

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$$(P^t f)(\mathbf{x}) = \lim_{M \rightarrow \infty} \mathbb{E} \tilde{f}_M(\mathbf{x}, \xi(t)),$$

whereas for the generator we have

$$(Af)(\mathbf{x}) = \lim_{M \rightarrow \infty} L_{\mathbf{y}} \tilde{f}_M(\mathbf{x}, \mathbf{y}) \Big|_{\mathbf{y}=0}.$$

Difference lives on the boundary

First lemma on the difference of semigroups

Let $f \in W_2^2(D)$ and $\mathbf{x} \in D$. Then

$$(P^t f)(\mathbf{x}) - (R^t f)(\mathbf{x}) = - \lim_{M \rightarrow \infty} \int_0^t P^\tau L(\tilde{f}_M - \bar{f}_M)(\mathbf{x}, 0) d\tau$$

and in $W_2^2(D)$.

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Remark. We shall employ this lemma to define a *pathwise* accumulated momentum.

Difference of semigroups lives on the boundary

Second lemma on the difference of semigroups

For every $f \in W_2^2(D)$ holds

$$(P^t f)(\mathbf{x}) - (R^t f)(\mathbf{x}) = \int_{\partial D} Q^t(\mathbf{x}, \mathbf{z})(\gamma_1 f)(\mathbf{z}) dS(\mathbf{z}),$$

where

$$Q^t(\mathbf{x}, \mathbf{z}) = \frac{1}{2} \int_0^t \tilde{R}^\tau(\mathbf{x}, \mathbf{z}) d\tau, \quad \tilde{R}^\tau(\mathbf{x}, \mathbf{z}) = \sum_{l=0}^{\infty} \frac{L(\kappa_l)}{\kappa_l^2} e^{-tL(\kappa_l)} s_l(\mathbf{x}) \overline{s_l(\mathbf{z})}.$$

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Remark. Recall that **local time is a measure on the boundary**.

It is exactly the case here: $d\mathcal{L}_\tau = \frac{1}{2} \tilde{R}^\tau(\mathbf{x}, \mathbf{z}) \delta_{\partial D}(\mathbf{z}) d\tau$.

A small digression regarding Brosmaler theorem

Let X_s be a reflecting BM in D and L_s its local time.

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Classical result due to Brosmaler [2]:

$$u(\mathbf{x}) = \frac{1}{2} \lim_{t \rightarrow \infty} \mathbb{E}_{\mathbf{x}} \int_0^t f(X_s) dL_s$$

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In our setting,

$$u(t, \mathbf{x}) = (Q^t g)(\mathbf{x})$$

solves $-Au = 0$, $u(0, \mathbf{x}) = 0$, $\gamma_1 u = g$.

Three corollaries on the group structure (1)

Theorem 1

The operator families $(R^t)_{t \geq 0}$ and $(Q^t)_{t \geq 0}$ satisfy the following evolution relations:

$$R^{t+s} = R^t R^s,$$

$$Q^{t+s} = Q^t + \tilde{R}^t Q^s$$

and $R^0 = I$, $Q^0 = 0$.

Three corollaries on the group structure (2)

Theorem 2

For every $t > 0$ and $f \in L_2(D)$ holds

$$\frac{\partial}{\partial t} R^t f = \frac{1}{2} A^N R^t f.$$

Three corollaries on the group structure (2)

Theorem 2

For every $t > 0$ and $f \in L_2(D)$ holds

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Theorem 3

For every $t > 0$ and $g \in W_2^{1/2}(\partial D)$ holds

$$\frac{\partial}{\partial t} Q^t g = \frac{1}{2} \int_{\partial D} \tilde{R}^t(\mathbf{x}, \mathbf{y}) g(\mathbf{y}) dS(\mathbf{y}).$$

Pathwise definition of the accumulated momentum

It is easy to guess a pathwise counterpart of P^t – it should be a shift by $\xi(t)$

$$(\mathcal{P}^t f)(\mathbf{x}) = f(\mathbf{x} + \xi(t))$$

(in fact it is preferable to choose something equivalent to this definition).

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(in fact it is preferable to choose something equivalent to this definition).

By way of analogy with the first lemma, define:

$$(\mathcal{Q}^t g)(\mathbf{x}) = \lim_{M \rightarrow \infty} \int_0^t \mathcal{P}^\tau L(\tilde{G}_M - \bar{G}_M)(\mathbf{x}, 0) d\tau, \quad (3)$$

where $G(\mathbf{x}, \mathbf{y})$ is a special continuation of g into the domain.

Properties of pathwise accumulated momentum

Existence

The limit in the rhs of (3) exists in $L_2(\mathcal{H}, \mu)$, where $\mathcal{H} = D \times \Omega$ and $d\mu = d\mathbf{x} \times d\mathbf{P}$.

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The average of Q^t

Let $g \in W_2^{1/2}(\partial D)$. Then

$$\mathbb{E}(Q^t g)(\mathbf{x}) = (Q^t g)(\mathbf{x}).$$

Limit theorem for the reflecting BM in d -ball

([5] by me). Let us approximate the BM $w(t)$ with a random walk

$$\zeta_n(t) = \frac{1}{\sqrt{n}} \sum_{j=1}^{\eta(nt)} \xi_j.$$

Here ξ_j are iidrvs with common rotation-invariant distribution and a unitary second moment.

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From this we derive the definition of Q_n^t .

Limit theorem for the reflecting BM in d -ball

Limit theorem for the process in the domain

Let $f \in D(A^N)$. Then

$$\|R_n^t f - R^t f\|_{L_2(D)} \leq \frac{C\sqrt{t}}{\sqrt{n}} \|f\|_{W_2^2(D)}.$$

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Limit theorem for the local time

For every $g \in W_2^{1/2}(\partial D)$ holds

$$\|Q_n^t g - Q^t g\|_{L_2(D)} \leq \frac{Ct^{3/8}}{n^{3/8}} \|g\|_{W_2^{1/2}(\partial D)}^2.$$

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