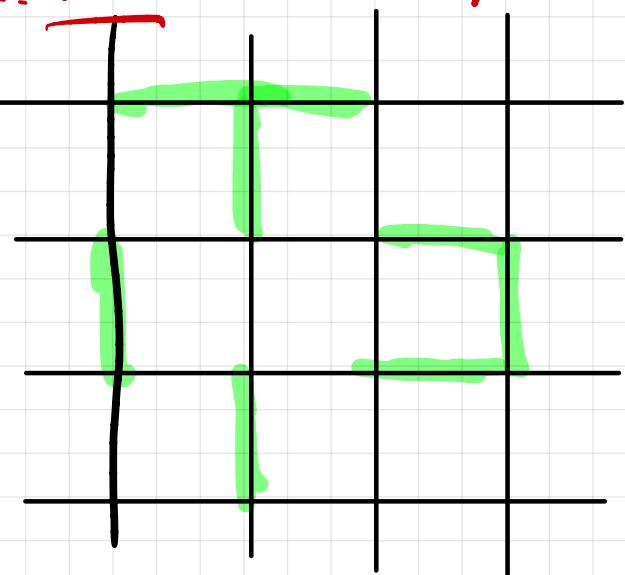


Lattice models of statistical mechanics

I. Def. Chapter I. Percolation



Hypercubic lattice:

- vertices:

$$\mathbb{Z}^d := \{(x_1, \dots, x_d) : x_i \in \mathbb{Z}\}$$

- edges:

$$E^d := \{(x, y) : |x - y|_1 = 1\}$$

$$= \{x - y = (0, \dots, 0, \pm 1, 0, \dots)\}$$

[x, y differ in one coordinate]

Let $G = (V, E) \subset \mathbb{Z}^d$ - finite subgraph.

Take $p \in [0, 1]$

Percolation model:

each edge is open with probability p , independently of the others;

$$\text{IP}(e \text{ is open}) = p$$

Def:

Space of configurations:

$$\Omega := \{\text{closed}, \text{open}\}^E$$

Closed

Open



$$\mathcal{N} = \{ (\omega_e)_{e \in E} : \omega_e \in \{0, 1\}\}$$

On every edge, we have
a Bernoulli rand. var.:

$$\begin{aligned} \omega_e &\xrightarrow{\quad} 1 \text{ (open)} \quad \text{w. prob. } p \\ &\xrightarrow{\quad} 0 \text{ (closed)} - u - 1-p. \end{aligned}$$

Probability measure on \mathcal{N} :
product of Bernoulli
distributions.

$$P_p(\omega) = p^{\# \text{ open } (w)} (1-p)^{\# \text{ closed } (w)}$$

Parameter p :

p ~ density of open edges

If $p \approx 0$, then almost no
open edges

If $p \approx 1$, then almost all
edges are open.

By the Law of Large Numbers,
the expected number of
open edges is $p \cdot |E|$.

Configurations \leftrightarrow subgraphs.

$\omega \in \{0, 1\}^E \hookrightarrow \Gamma = (V, \text{edges } w_e = 1)$

Subgraph \uparrow or open edges.

Def (on (\mathbb{R}^d, E^d))

$\mathcal{N} = \{0, 1\}^{E^d}$

6-algebra \mathcal{F} : generated by cylinder events

[events that depend on finitely many edges]

Percolation measure on \mathcal{N}

is the product of Bernoulli rand. var. on edges $e \in E^d$

[Kolmogorov's extension thm.
Carathéodory's thm.]

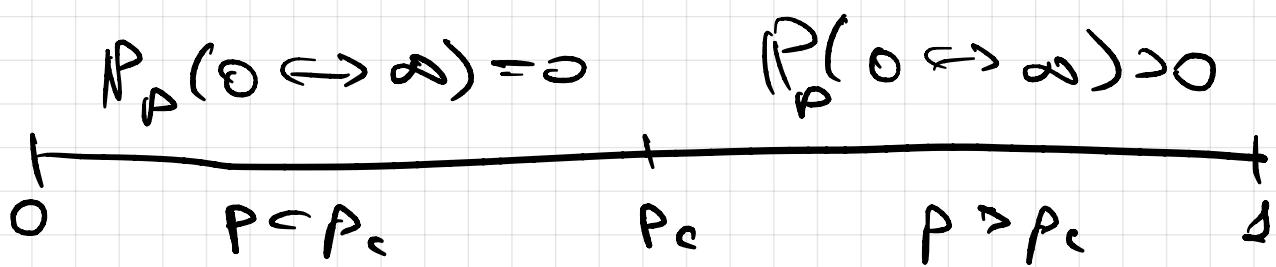
Basic question:

is there an infinite connected component
of open edges?

Or that the component of \mathcal{O}
is infinite? $\mathcal{O} \xrightarrow{\quad} \mathcal{O}$

③

Depends on P !



Monotonicity in P ?

2. Monotonicity in P

Notations:

Let $x, y \in \mathbb{Z}^d$. Events:

- $\{x \leftrightarrow y\} := \{\exists \text{ path on open edges from } x \text{ to } y\}$
- $\{x \leftrightarrow \infty\} := \{x \text{ belongs to the infinite component } y \text{ (cluster)}\}$

ExFC

$\{x \leftrightarrow y\}, \{x \leftrightarrow \infty\}$ are in the σ -algebra \mathcal{F} .

Def:

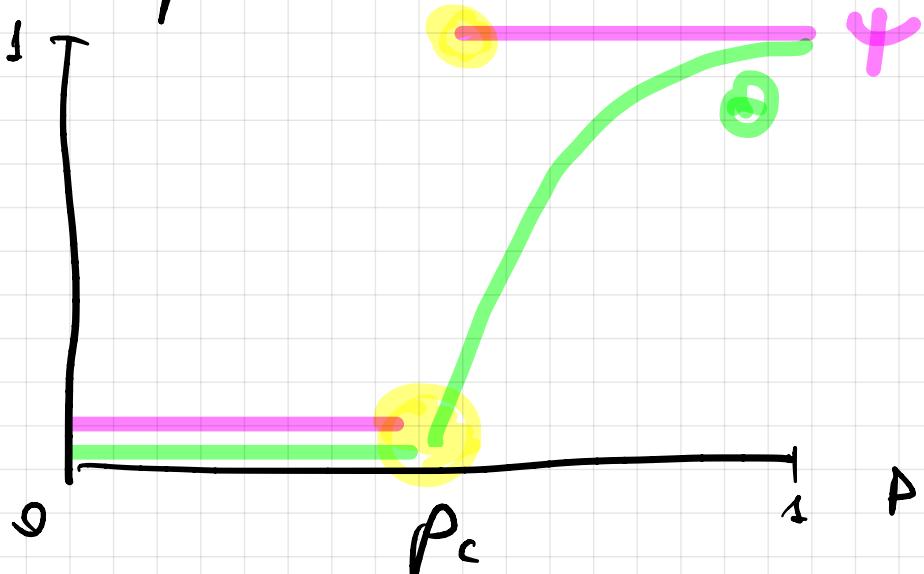
- Connectivity function:

$$\Theta(P) := P_P(0 \leftrightarrow \infty)$$

- Existence of an infinite cluster:

$$\Upsilon(P) := P_P\left(\bigcup_{x \in \mathbb{Z}^d} \{x \leftrightarrow \infty\}\right)$$

Expected behaviour:



Big open problem:

$$\Theta(p_c) = 0 \text{ in } \mathbb{Z}^d.$$

Prop. (monotonicity)

ψ, θ are non-decreasing in p .

Proof:

We'll construct a coupling for percolation measures with different p .

On every edge, we put a uniform rand. var. U_2 on $[0, 1]$,

that are indep. at each other!

$\{U_e\}_{e \in E^d}$: $U_e \sim \text{uniform}[0, 1]$
 U_e and U_f are indep.
 if $e \neq f$.

Let $p \in [0, 1]$.

Define ω (percolation config.):

- $\omega_e = 1$ if $U_e \leq p$
- $\omega_e = 0$ if $U_e > p$

Claim: $\omega \sim IP_p$

Proof:

$$IP(U_e \leq p) = p.$$

$$\text{So } IP(\omega_e = 1) = p.$$

$$\text{and } IP(\omega_e = 0) = 1 - p.$$

Every ω_e is a function
of U_e .

Since $(U_e)_{e \in E^d}$ are i.i.d.,

so are $(\omega_e)_{e \in E^d}$



Let $q \in (p, 1)$.

Define ω' :

- $\omega'_e = 1$ if $U_e \leq q$
- $\omega'_e = 0$ if $U_e > q$.

As before, $\omega' \sim \text{IP}_q$.

Consequently:

$w_e \leq w'_e$ at every $e \in E$,

Indeed, if $w_e = 1$, then
 $u_e \leq p \leq q$. Hence $w'_e = 1$.

Then, if $\omega \notin \{0 \leftrightarrow \infty\}$,
then also $\omega' \notin \{0 \leftrightarrow \infty\}$.

Hence,

$$\text{IP}(\overset{\omega}{0 \leftrightarrow \infty}) \leq \text{IP}(\overset{\omega'}{0 \leftrightarrow \infty})$$

" " "

$$\Theta(p) = \text{IP}_p(0 \leftrightarrow \infty)$$

$$\text{IP}_q(0 \leftrightarrow \infty) = \Theta(q)$$

For γ - analogously

3. Phase transition

Def:

Critical density:

$$p_c := \inf \{p : \Theta(p) > 0\}$$

Thm(i) $p_c(\infty) = 1$. [no phase transition](ii) $d \geq 2$, then

$$0 < p_c(\infty^d) < 1$$

Proof:(ii) Part 1: $p_c > 0$.

We want to prove that when p is small enough,

$$\lim_{n \rightarrow \infty} P_p(0 \leftrightarrow \infty) = 0.$$

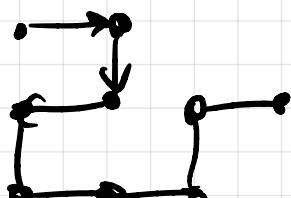
We will use a counting argument:

The weight of a long path will dominate the number of such paths.

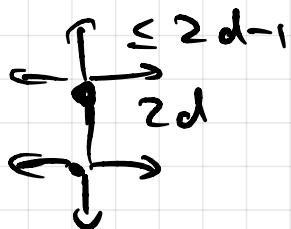
Paths: (simple)

self-avoiding walks

SAW_n : set of all self-avoiding walks of length n , starting from 0 ,



$P_p(\emptyset \hookrightarrow \infty) = P_p(\text{t}_n \exists \delta \in SAW_n : \text{all edges in } t \text{ are open})$



$$\begin{aligned}
 & P_p(\emptyset \hookrightarrow \infty) \leq \sum_{\delta \in SAW_n} P_p(\delta \text{ is open}) \\
 & = p^n \cdot |SAW_n| \\
 & \leq p^n \cdot 2d \cdot (2d-1)^{n-1} \\
 & \quad \text{if } p < \frac{1}{2d-1} \rightarrow 0
 \end{aligned}$$

Thus, for any $p < \frac{1}{2d-1}$:

$$P_p(\emptyset \hookrightarrow \infty) = 0.$$

$$\text{Hence } P_c \geq \frac{1}{2d-1}.$$